

Modified Method of Characteristics Variational Multiscale Finite Element Method for Time Dependent Navier-Stokes Problems*

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Abstract. In this paper, a modified method of characteristics variational multiscale (MMOCVMS) finite element method is presented for the time dependent Navier-Stokes problems, which is leaded by combining the characteristics time discretization with the variational multiscale (VMS) finite element method in space. The theoretical analysis shows that this method has a good convergence property. In order to show the efficiency of the MMOCVMS finite element method, some numerical results of analytical solution problems are presented. First, we give some numerical results of lid-driven cavity flow with $Re = 5000$ and 7500 as the time is sufficient long. From the numerical results, we can see that the steady state numerical solutions of the time-dependent Navier-Stokes equations are obtained. Then, we choose $Re = 10000$, and we find that the steady state numerical solution is not stable from $t = 200$ to 300 . Moreover, we also investigate numerically the flow around a cylinder problems. The numerical results show that our method is highly efficient.

Keywords: modified method of characteristics, VMS finite element method, time dependent Navier-Stokes problems, characteristics-based method, numerical tests.

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1 Introduction

In this paper, we consider the time-dependent Navier-Stokes problems

$$\begin{cases} u_t - \nu \Delta u + (u \cdot \nabla)u + \nabla p = f, & x \in \Omega \times [0, T], \\ \nabla \cdot u = 0, & x \in \Omega \times [0, T], \\ u(x, 0) = u_0(x), & x \in \Omega, \\ u(x, t) = 0, & x \in \partial\Omega \times [0, T], \end{cases} \quad (1.1)$$

where Ω is a bounded domain in \mathbb{R}^2 assumed to have a Lipschitz continuous boundary $\partial\Omega$. $u = (u_1(x, t), u_2(x, t))^T$ represents the velocity vector, $p(x, t)$ represents the pressure, $f(x, t)$ is the body force, $\nu = 1/Re$ the viscosity number, Re is the Reynolds number.

Developing efficient finite element methods for the Stokes and Navier-Stokes equations is a key component in the incompressible flow simulation. There are some iterative methods for solving the stationary Navier-Stokes equations under some strong uniqueness conditions were presented by some authors [15, 18, 19]. The VMS method was first introduced by Hughes and his coworkers in [22, 23]. The basic idea is splitting the solution into resolved and unresolved scale, representing the unresolved scales in terms of the resolved scales, and using this representation in the variational equation for the resolved scales. And there are many works devoted to this method, e.g. VMS method for the Navier-Stokes equations [26]; a two-level VMS method for convection-dominated convection-diffusion problems [29]; VMS methods for turbulent flow [9, 30, 31]; large-eddy simulation (LES) [24, 25, 32]; subgrid-scale models for the incompressible flow [22, 44]. In [27, 28], John et al. presented the error analysis of the two different kinds of VMS for the Navier-Stokes equations. The main difference is the definition of the large scales projection (L^2 -projection in [28] and elliptic projection in [27]). There is another class of VMS method which rely on a three-scale decomposition of the flow field into large, resolved small and unresolved scales [8].

In the characteristics method, which is a highly effective method for advection dominated problems, the hyperbolic part (the temporal and advection term) is treated by a characteristic tracking scheme. In 1982, Douglas and Russell [11] presented the modified method of characteristics finite element method (MMOCFEM) firstly. Russell [36] extended it to nonlinear coupled systems in two and three spatial dimensions. The second order in time method for linear convection diffusion problems had been given by Ewing and Russell [13]. A characteristics mixed finite element method for advection-dominated transport problems was presented by Arbogast [2] and for Navier-stokes equations by Buscaglia and Dari [7], respectively. In [34], a detail theory analysis for Navier-Stokes equations have been done by Pironneau, who obtained suboptimal convergence rates of the form $\mathcal{O}(h^m + \Delta t + h^{m+1}/\Delta t)$ and improved by Dawson et al [10]. In [4], Boukir et al. presented a second-order time scheme based on the characteristics method and spatial discretization of finite element type for the incompressible Navier-Stokes equations. An optimal error estimate for the Lagrange-Galerkin mixed finite element approximation of the Navier-Stokes equations was given by Süli in [43]. In [38, 39], one order and

second order MMOC mixed defect-correction finite element methods for time-dependent Navier-Stokes problems were given. In [41], second order in time MMOC variational multiscale finite element method for the time-dependent Navier-Stokes equation was shown. In [40], we presented modified characteristics gauge-Uzawa finite element method for the conduction-convection equation. In [42], the modified characteristics finite element method for the Navier-Stokes/Darcy problem was shown.

In this paper, we present a modified method of characteristics VMS finite element method based on the L^2 projection for the time dependent Navier-Stokes equations. In our method, the hyperbolic part (the temporal and advection term) is treated by a characteristic tracking scheme. The VMS based on projection finite element method is used for spatial discretization. The error analysis shows that this method has a good convergence property. In order to show the efficiency of the MMOCVMS finite element method, we firstly present some numerical results of analytical solution problems. The numerical results show that the convergence rates are $\mathcal{O}(h^3)$ in the L^2 -norm for u , $\mathcal{O}(h^2)$ in the semi H^1 -norm for u and $\mathcal{O}(h^2)$ in the L^2 -norm for p by using the Taylor-Hood element, which agrees very well with our theoretical results. Then, some numerical results of the lid-driven cavity flow with $Re = 5000$ and 7500 are given. We present the numerical results as the time are sufficient long enough, then the solution of the Navier-Stokes should approximate the solution of the steady state cases. From the numerical results, we can see that a steady state numerical solutions of the time-dependent Navier-Stokes equations are obtained. And the numerical solutions are in good agreement with that of the steady Navier-Stokes equations shown by Ghia et al. [14] and Erturk et al. [12]. Finally, we present some numerical results for $Re = 10000$. It shows that the solution is mainly periodic with small variations in the amplitude of the time evolution at the monitoring points as the time is sufficient long. And the phase portraits of the monitoring points show that the variations in amplitude yield a solution which is quasi-periodic. Furthermore, we present numerical simulations for the two-dimensional fluid flow around a cylinder. It is observed from these numerical results that the scheme can result in good accuracy, which shows that this method is highly efficient.

The rest of this paper is organized as follows. In Section 2, the functional settings of the Navier-Stokes equations are presented. In Section 3 the MMOCVMS finite element methods is proposed. The following Section 4 we derive optimal error estimate for the new algorithm. Numerical experiments will be carried out in Section 5 and some concluding remarks will be given in the final section.

2 Functional settings of the Navier-Stokes equations

In this section, we aim to describe some notations and results which will be frequently used in this paper. Firstly we introduce the standard Sobolev spaces

$$H^n(0, T; \Phi) = \left\{ v \in \Phi; \sum_{0 \leq i \leq n} \int_0^T \left(\left\| \frac{\partial^i v}{\partial t^i} \right\|_{\Phi} \right)^2 dt < +\infty \right\},$$

equipped with the norm

$$\|v\|_{H^n(\Phi)} = \left[\sum_{0 \leq i \leq n} \int_0^T \left(\left\| \frac{\partial^i v}{\partial t^i} \right\|_{\Phi} \right)^2 dt \right]^{\frac{1}{2}}.$$

Here Φ denotes a Hilbert space with the norm $\|\cdot\|_{\Phi}$. Especially, when $n = 0$ we note

$$\|v\|_{L^2(\Phi)} = \left(\int_0^T \|v\|_{\Phi}^2 dt \right)^{\frac{1}{2}}.$$

And, we define

$$L^{\infty}(0, T; \Phi) = \left\{ v \in \Phi; \operatorname{ess} \sup_{0 \leq t \leq T} \|v\|_{\Phi} < +\infty \right\}$$

with the norm

$$\|v\|_{L^{\infty}(\Phi)} = \operatorname{ess} \sup_{0 \leq t \leq T} \|v\|_{\Phi}.$$

For the mathematical setting of the Navier-Stokes problems (1.1), we introduce the following Hilbert spaces

$$\begin{aligned} X &:= H_0^1(\Omega)^2, \quad M := L_0^2(\Omega) = \left\{ \varphi \in L^2(\Omega); \int_{\Omega} \varphi dx = 0 \right\}, \\ V &:= \{v \in X; (\nabla \cdot v, q) = 0 \forall q \in M\}. \end{aligned}$$

The following assumptions and results are recalled (see [18]).

(A₁) There exists a positive constant C_0 which only depends on Ω , such that

$$\|u\|_0 \leq C_0 \|\nabla u\|_0 \quad \forall u \in X.$$

(A₂) Assume that Ω is smooth, hence the unique solution $(v, q) \in (X, M)$ of the steady Stokes problem

$$-\Delta v + \nabla q = g, \quad \nabla \cdot v = 0, \quad \text{in } \Omega, \quad v|_{\partial\Omega} = 0$$

for any prescribed $g \in L^2(\Omega)^2$ exists and satisfies

$$\|v\|_2 + \|q\|_1 \leq c \|g\|_0,$$

where $c > 0$ is a generic constant depending on Ω , which may stand for different values at its different occurrences. $\|\cdot\|_2$ means the H^2 -norm, $\|\cdot\|_1$ means the H^1 -norm, and $\|\cdot\|_0$ means the L^2 -norm, respectively.

For simplicity, we assume that the time-dependent Navier-Stokes system (1.1) has a unique solution (u, p) satisfying

$$\begin{aligned} &\|u_0\|_{H^{r+1}} + \|u\|_{L^{\infty}(I; H^{r+1})} + \|u_t\|_{L^2(I; H^{r+1})} + \|u_{tt}\|_{L^2(I; L^2)} \\ &\quad + \|p\|_{L^{\infty}(I; H^r)} + \|p_t\|_{L^2(I; H^r)} \leq C \end{aligned}$$

with $I = (0, T)$ and $r \geq 2$.

3 MMOCVMS finite element methods

3.1 MMOC-mixed finite element methods for the time-dependent Navier-Stokes equations

For each positive integer N , let $\{\mathcal{J}_n : 1 \leq n \leq N\}$ be a partition of $[0, T]$ into subintervals $\mathcal{J}_n = (t_{n-1}, t_n]$, with $t_n = n\Delta t$, $\Delta t = T/N$. Set $u^n = u(\cdot, t_n)$ be the $u(x, t)$ at $t = t_n$. The characteristic trace-back along the field u^{n-1} of a point $x \in \Omega$ at time t_n to t_{n-1} is approximately $\bar{x}(x, t_{n-1}) = x - \Delta t u^{n-1}$.

Consequently, the hyperbolic part in the first equation of (1.1) at time t_n is approximated by

$$u_t + u^{n-1} \cdot \nabla u^n \approx \frac{u^n - \bar{u}^{n-1}}{\Delta t},$$

where

$$\bar{w} = \begin{cases} w(\bar{x}), & \bar{x} = x - \Delta t u^{n-1} \in \Omega, \\ 0, & \text{otherwise} \end{cases}$$

for any function w .

Let \mathfrak{S}_h be a quasi-uniform partition of $\bar{\Omega}$ into non-overlapping triangles, indexed by a parameter $h = \max_{K \in \mathfrak{S}_h} \{h_K; h_K = \text{diam}(K)\}$. We introduce the finite element subspace $X_h \subset X$, $M_h \subset M$ as follows

$$\begin{aligned} X_h &= \{v_h \in X \cap C^0(\bar{\Omega})^2; v_h|_K \in P_\ell(K)^2 \forall K \in \mathfrak{S}_h\}, \\ M_h &= \{q_h \in M \cap C^0(\bar{\Omega}); q_h|_K \in P_k(K) \forall K \in \mathfrak{S}_h\}, \end{aligned}$$

where $P_\ell(K)$ is the space of polynomials of degree ℓ on K . $\ell \geq 1, k \geq 1$ are two integers. Assume that (X_h, M_h) satisfies the discrete LBB condition

$$\sup_{v_h \in X_h} \frac{d(\varphi_h, v_h)}{\|\nabla v_h\|_0} \geq \beta \|\varphi_h\|_0 \quad \forall \varphi_h \in M_h,$$

where $d(\varphi, v) = (\varphi, \nabla \cdot v)$, $\beta > 0$ is a constant independent of h . In this paper, V_h denotes the kernel of the discrete divergence operator

$$V_h = \{v_h \in X_h; (q_h, \nabla \cdot v_h) = 0 \quad \forall q_h \in M_h\}.$$

With the previous notations, we present the MMOC finite element method for the Navier-Stokes (1.1): find $(u^n, p^n) : \{t_1, \dots, t_N\} \rightarrow V_h \times M_h$ such that

$$\begin{aligned} \left(\frac{u_h^n - \tilde{u}_h^{n-1}}{\Delta t}, v_h \right) + \nu a(u_h^n, v_h) - d(p_h^n, v_h) &+ d(\varphi_h, u_h^n) = (f^n, v_h) \\ &\forall v_h \in X_h, \varphi_h \in M_h, \end{aligned} \quad (3.1)$$

where $a(u, v) = (\nabla u, \nabla v)$, $f^n = f(t_n)$,

$$\tilde{u}_h^{n-1} = \begin{cases} u_h^{n-1}(\tilde{x}), & \tilde{x} = x - u_h^{n-1} \Delta t \in \Omega, \\ 0, & \text{otherwise.} \end{cases}$$

3.2 The VMS finite element method and MMOCVMS finite element method for the time-dependent Navier-Stokes equations

For the VMS finite element method, a large scale space $L_H \subset L = \{\mathbb{L} \in (L^2(\Omega))^{d \times d}, \mathbb{L} = \mathbb{L}^T\}$ and a so-called turbulent viscosity $\nu_T \geq 0$ are introduced. The semi-discrete problems reads as follows (please see [31]): Find $u_h : [0, T] \rightarrow v_h, p_h : [0, T] \rightarrow M_h$ and $\mathbb{G}_H : [0, T] \rightarrow L_H$ satisfying

$$\begin{aligned} (u_{ht}, v_h) + (\nu + \nu_T)a(u_h, v_h) - d(p_h, v_h) + d(\varphi_h, u_h) + b(u_h, u_h, v_h) \\ - (\nu_T \mathbb{G}_H, \nabla v_h) = (f, v_h) \quad \forall v_h \in X_h, \varphi_h \in M_h, \\ (\mathbb{G}_H - \nabla u_h, \mathbb{L}_h) = 0 \quad \forall \mathbb{L}_h \in L_H, \end{aligned} \quad (3.2)$$

where $u_h(x, 0) = u_h^0 \in V_h$ is a discretely divergence free approximation of u_0 ,

$$b(u, v, w) = \frac{1}{2} \left[\int_{\Omega} \sum_{i,k=1}^2 u_i \frac{\partial v_k}{\partial x_i} w_k - \sum_{i,k=1}^2 u_i \frac{\partial w_k}{\partial x_i} v_k dx \right] \quad \forall u, v, w \in X_h.$$

Let $P_{L_H} : L \rightarrow L_H$ for all $v \in V$, with

$$(P_{L_H} \nabla v - \nabla v, \mathbb{L}_H) = 0 \quad \forall \mathbb{L}_H \in L_H$$

denote the L^2 -projection from L onto L_H . Then, $\mathbb{G}_H = P_{L_H} \nabla u_h$ in (3.2). Since P_{L_H} is an L^2 -projection, it follows for $v \in V$ and $\|\nabla v\|_0 > 0$,

$$\begin{aligned} \nu_T \|(I - P_{L_H}) \nabla v\|_0^2 &= \nu_T (\|\nabla v\|_0^2 - \|P_{L_H} \nabla v\|_0^2) \\ &= \nu_T \left(1 - \frac{\|P_{L_H} \nabla v\|_0^2}{\|\nabla v\|_0^2} \right) \|\nabla v\|_0^2 \equiv \nu_{add} \|\nabla v\|_0^2. \end{aligned} \quad (3.3)$$

By $0 \leq \|P_{L_H} \nabla v\|_0 \leq \|\nabla v\|_0$, we get $0 \leq \nu_{add} \leq \nu_T$. By straightforward calculation, we deduce

$$(\nu_T \nabla u_h, \nabla v_h) - (\nu_T P_{L_H} \nabla u_h, \nabla v_h) = (\nu_T (I - P_{L_H}) \nabla u_h, (I - P_{L_H}) \nabla v_h). \quad (3.4)$$

Hence, system (3.2) can be reformulated as: Find $u_h : [0, T] \rightarrow X_h, p_h : [0, T] \rightarrow M_h$ satisfying

$$\begin{aligned} (u_{ht}, v_h) + \nu a(u_h, v_h) &- d(p_h, v_h) + d(\varphi_h, u_h) + b(u_h, u_h, v_h) \\ &+ (\nu_T (I - P_{L_H}) \nabla u_h, (I - P_{L_H}) \nabla v_h) \\ &= (f, v_h) \quad \forall v_h \in X_h, \varphi_h \in M_h. \end{aligned}$$

The MMOC time discretization, combined with the VMS finite element method in space, leads to the following MMOCVMS finite element method.

Find $(u_h^n, p_h^n) \in (X_h \cap V_h) \times M_h$, such that

$$\begin{aligned} \left(\frac{u_h^n - \hat{u}_h^{n-1}}{\Delta t}, v_h \right) + (\nu + \nu_T)a(u_h^n, v_h) - d(p_h^n, v_h) + d(\varphi_h, u_h^n) \\ - (\nu_T \mathbb{G}_H^n, \nabla v_h) = (f^n, v_h) \quad \forall v_h \in X_h, \varphi_h \in M_h, \\ (\mathbb{G}_H^n - \nabla u_h^n, \mathbb{L}_H) = 0 \quad \forall \mathbb{L}_H \in L_H, \end{aligned} \quad (3.5)$$

where

$$\hat{u}_h^{n-1} = \begin{cases} u_h^{n-1}(\hat{x}), & \hat{x} = x - u_h^{n-1}\tau \in \Omega, \\ 0, & \text{otherwise.} \end{cases}$$

Using (3.4), (3.5) can be reformulated as

$$\begin{aligned} & \left(\frac{u_h^n - \hat{u}_h^{n-1}}{\Delta t}, v_h \right) + \nu a(u_h^n, v_h) - d(p_h^n, v_h) + d(\varphi_h, u_h^n) \\ & + (\nu_T(I - P_{L_H})\nabla u_h, (I - P_{L_H})\nabla v_h) = (f^n, v_h) \quad \forall v_h \in X_h, \varphi_h \in M_h. \end{aligned} \quad (3.6)$$

Remark 1. Since L_H has been distinguish between resolved small scales and large scales, with L_H representing the large scales, it must be in some sense a coarse finite element space. One way of achieving this is choosing it to be a lower order finite element space than V_h on the same grid (e.g. $L_H = \mathbb{D}(V_h)$, $\mathbb{D}(v) = (\nabla v + \nabla v^T)/2$, see [27] for details). The other way is defining L_H on a coarser grid. This two methods are called one-level and two-level projection-based FEVMS methods respectively (see [29, 31]). The so-called turbulent viscosity ν_T can be chosen in many ways. In this paper we choose as follows (Smagorinsky-type [26]) $\nu_T = \delta h^2 \|\nabla u_h^n\|_0$, where δ is a suitable constant.

Remark 2. [35, 44] If $L_H = \{v_h \in (C^0(\Omega))^2 \cap X : v_h|_K \in (P_k(K))^2 \forall K \in \mathfrak{S}_h\}$, then we have the following properties

$$\begin{aligned} \|P_{L_H} v\|_1 & \leq C \|v\|_1 \quad \forall v \in X, \quad \|\nabla(I - P_{L_H})v\|_0 \leq C \|\nabla v\|_0 \quad \forall v \in X, \\ \|v - P_{L_H} v\|_1 & \leq Ch^l \|v\|_{l+1}, \quad 1 \leq l \leq k \quad \forall v \in H^{l+1}(\Omega)^2 \cap X. \end{aligned}$$

4 Error estimate

In order to obtain the error analysis, we give some lemmas firstly.

Lemma 1. [39] *Let*

$$e(x, n) = \frac{u^n(x) - \bar{u}^{n-1}(x)}{\Delta t} - \left(\frac{\partial u}{\partial t}(x, t_n) + u^n(x) \nabla u^n(x) \right)$$

and $\tau > 0$ such that $u \in \mathcal{C}^4([\tau, T]; H^3(\Omega)^2)$. For $t_n > \tau$, we have

$$e(x, n) = -\Delta t \left(\frac{1}{2} \frac{d^2 g_x^n}{dt^2} + \frac{\partial u}{\partial t} \cdot \nabla u(x, t_n) \right) + O(\Delta t^2), \quad (4.1)$$

where $g_x^n(t) = u(x - (t_n - t)u^{n-1}, t)$, $u^n(x) = u(x, t_n)$.

Then, we define the following Galerkin projection

$$(R_h, Q_h) = (R_h(u, p), Q_h(u, p)) : (X, M) \rightarrow (X_h, M_h),$$

such that

$$\begin{aligned} \nu a(R_h - u, v_h) & - d(Q_h - p, v_h) + d(\varphi_h, R_h - u) = 0 \\ & \forall (u, p) \in (X, M), (v_h, \varphi_h) \in (X_h, M_h). \end{aligned} \quad (4.2)$$

Lemma 2. [20, 33] The Galerkin projection (R_h, Q_h) satisfies

$$\begin{aligned} \|R_h - u\|_0 + h(\|\nabla(R_h - u)\|_0 + \|Q_h - p\|_0) \\ \leq Ch^{k+1}(\nu\|u\|_{k+1} + \|p\|_k), \quad k = 1, 2. \end{aligned} \quad (4.3)$$

Lemma 3. [1, 3] The Galerkin projection (R_h, Q_h) satisfies

$$\|R_h\|_\infty \leq C(\|u\|_\infty + h^k\|u\|_{H^{k+1}(\Omega)}).$$

4.1 Error estimate for the velocity

Lemma 4. Let u_h^n be defined by (3.6). If $\Delta t \leq \min\{1/(2L_n), h\}$, $L_n = \max_{1 \leq l \leq n-1} \|\nabla u_h^l\|_\infty$, for all $2 \leq n \leq N$ we have

$$\begin{aligned} \|\xi_h^n\|_0^2 + \nu\Delta t \sum_{i=1}^n \|\nabla \xi_h^i\|_0^2 \\ \leq C \exp(CT) (\Delta t^2 + h^{2k+2} + \nu_{add}h^{2k} + \nu_T\|(I - P_{L_H})\nabla u^n\|_0^2), \end{aligned} \quad (4.4)$$

$$\|u_h^n\|_\infty < +\infty, \quad \Delta t\|u_h^n\|_{W^{1,\infty}} \leq 1/4, \quad (4.5)$$

where $\xi_h^n = u_h^n - R_h^n$, C is a positive constant independent of Δt and h .

Remark 3. Define $\hat{\mathcal{X}}_x^n(t) = x - (t_n - t)u_h^{n-1} \forall t \in [t^{n-1}, t^n]$, $2 \leq n \leq N$. Since X_h is a subset of $W^{1,\infty}(\Omega)$, under the condition $\Delta t \leq \frac{1}{2L_n}$ on the time step it is an easy matter to verify that this mapping has a positive Jacobian, since u_h^l vanishes on $\partial\Omega$; this mapping is one-to-one and this it is a change of variables from Ω onto Ω . This yields for any positive function ϕ on Ω the estimate (please see [4] for details)

$$\int_\Omega \phi(\hat{\mathcal{X}}_h^n(t))dx \leq C \int_\Omega \phi(x)dx.$$

Remark 4. Here, we assume that the initial condition u_0 has a strong regularity. Or there are some restriction on the time step Δt [17].

Proof. We prove this lemma by induction. By the definition of \hat{x} , we can see that (4.5) and (4.4) hold for $n = 0$. We assume that (4.5) and (4.4) hold for $1 \leq n \leq l-1$, then $L_n < +\infty$. Now we will prove them hold for $n = l$.

Letting $\varphi_h = 0$ in (3.6) and using the definition of V_h , we have

$$\begin{aligned} \left(\frac{u_h^l - \hat{u}_h^{l-1}}{\Delta t}, v_h \right) + \nu a(u_h^l, v_h) \\ + \nu_T((I - P_{L_H})\nabla u_h^l, (I - P_{L_H})\nabla v_h) = (f^l, v_h) \quad \forall v_h \in V_h. \end{aligned} \quad (4.6)$$

Subtracting $\left((R_h^l - \hat{R}_h^{l-1})/\Delta t, v_h \right) + \nu a(R_h^l, v_h)$ from both sides of (4.6), we get

$$\begin{aligned} \left(\frac{(u_h^l - R_h^l) - (\hat{u}_h^{l-1} - \hat{R}_h^{l-1})}{\Delta t}, v_h \right) + \nu a(u_h^l - R_h^l, v_h) \\ + \nu_T((I - P_{L_H})\nabla u_h^l, (I - P_{L_H})\nabla v_h) \\ = (f^l, v_h) - \left(\frac{R_h^l - \hat{R}_h^{l-1}}{\Delta t}, v_h \right) - \nu a(R_h^l, v_h). \end{aligned} \quad (4.7)$$

Define $\eta^l = u^l - R_h^l$, we arrive at

$$\begin{aligned}
 & \left(\frac{\xi_h^l - \xi_h^{l-1}}{\Delta t}, v_h \right) + \nu a(\xi_h^l, v_h) + \nu_T((I - P_{L_H})\nabla \xi_h^l, (I - P_{L_H})\nabla v_h) \\
 &= - \left(\frac{u^l - \bar{u}^{l-1}}{\Delta t} - \nu \Delta u^l + \nabla p^l - f^l, v_h \right) + \left(\frac{\hat{u}^{l-1} - \bar{u}^{l-1}}{\Delta t}, v_h \right) \\
 &+ \left(\frac{\eta^l - \hat{\eta}^{l-1}}{\Delta t}, v_h \right) + \left(\frac{\hat{\xi}_h^{l-1} - \xi_h^{l-1}}{\Delta t}, v_h \right) \\
 &+ \nu_T((I - P_{L_H})\nabla \eta_h^l, (I - P_{L_H})\nabla v_h) + \nu_T((I - P_{L_H})\nabla u^l, (I - P_{L_H})\nabla v_h) \\
 &+ (\nabla p^l, v_h) + \nu a(u^l - R_h^l, v_h) \tag{4.8}
 \end{aligned}$$

$$= \sum_{i=1}^6 \Pi_i + d(Q_h^l - p^l, v_h) + \nu a(u^l - R_h^l, v_h) = \sum_{i=1}^6 \Pi_i, \tag{4.9}$$

where Π_i denotes the i -th term of the right-hand side of (4.8).

Next we give the estimate for each term. By the definition of \hat{x} and \bar{x} , we arrive at

$$\hat{x}(x, t_{l-1}) - \bar{x}(x, t_{l-1}) = (u_h^{l-1} - u^{l-1})\Delta t.$$

Using Taylor formula, we obtain

$$\begin{aligned}
 \|\hat{u}^{l-1} - \bar{u}^{l-1}\|_0 &= \|u^{l-1}(\hat{x}) - u^{l-1}(\bar{x})\|_0 \\
 &\leq \Delta t \|\nabla u^{l-1}\|_\infty \|u_h^{l-1} - u^{l-1}\|_0 \\
 &\leq \Delta t \|\nabla u^{l-1}\|_\infty (\|u_h^{l-1} - R_h^{l-1}\|_0 + \|R_h^{l-1} - u^{l-1}\|_0).
 \end{aligned}$$

Therefore, we get

$$\begin{aligned}
 \frac{1}{\Delta t} \|\hat{u}^{l-1} - \bar{u}^{l-1}\|_0 &\leq \|\nabla u^{l-1}\|_\infty (\|u^{l-1} - R_h^{l-1}\|_0 + \|R_h^{l-1} - u^{l-1}\|_0) \\
 &\leq C(h^{k+1} + \|\xi_h^{l-1}\|_0). \tag{4.10}
 \end{aligned}$$

Now, we estimate the bound of $\|\frac{\eta^l - \hat{\eta}^{l-1}}{\Delta t}\|_0$. By Cauchy-Schwarz inequality, we get

$$\begin{aligned}
 \|\eta^l - \eta^{l-1}\|_0 &= \left(\int_\Omega (\eta^l - \eta^{l-1})^2 dx \right)^{\frac{1}{2}} = \left(\int_\Omega \left| \int_{t_{l-1}}^{t_l} \frac{\partial \eta}{\partial t}(x, \theta) d\theta \right|^2 dx \right)^{\frac{1}{2}} \\
 &\leq \sqrt{\Delta t} \left(\int_\Omega \int_{t_{l-1}}^{t_l} \left| \frac{\partial \eta}{\partial t}(x, \theta) \right|^2 d\theta dx \right)^{\frac{1}{2}} \leq \sqrt{\Delta t} \left\| \frac{\partial \eta}{\partial t} \right\|_{L^2([t_{l-1}, t_l]; L^2(\Omega))}. \tag{4.11}
 \end{aligned}$$

By the definition of $\hat{\mathcal{X}}_x^l(t_{l-1})$, we have

$$J(\hat{\mathcal{X}}_x^l(t_{l-1})) = \begin{pmatrix} 1 - \partial_x u_{h1}^{l-1} \Delta t & -\partial_y u_{h1}^{l-1} \Delta t \\ -\partial_x u_{h2}^{l-1} \Delta t & 1 - \partial_y u_{h2}^{l-1} \Delta t \end{pmatrix}. \tag{4.12}$$

Hence, $\det J(\hat{\mathcal{X}}_x^l(t_{l-1})) = 1 + \mathcal{O}(\Delta t)$. Then, we deduce

$$\begin{aligned} \|\eta^{l-1} - \hat{\eta}^{l-1}\|_{-1} &= \sup_{v \in V} (\|\nabla v\|_0^{-1} (\eta^{l-1} - \hat{\eta}^{l-1}, v)) = \sup_{v \in V} \left[\|\nabla v\|_0^{-1} \right. \\ &\quad \times \left. \left(\int_{\Omega} \eta^{l-1}(x) v(x) dx - \int_{\Omega} \eta^{l-1}(z) v(\hat{\mathcal{X}}_x^n(t_{l-1})^{-1}) (1 + \mathcal{O}(\Delta t^2)) dz \right) \right] \\ &\leq \sup_{v \in V} \left(\|\nabla v\|_0^{-1} \int_{\Omega} \eta^{l-1}(x) (v(x) - v(\hat{\mathcal{X}}_x^l(t_{l-1})^{-1})) dx \right) \\ &\quad + \sup_{v \in V} \left(C \Delta t^2 \|\nabla v\|_0^{-1} \int_{\Omega} \eta^{l-1}(z) v(\hat{\mathcal{X}}_x^l(t_{l-1})^{-1}) dz \right). \end{aligned}$$

Let $G(x) = x - \hat{\mathcal{X}}_x^l(t_{l-1})^{-1}$, then $|G(x)| \leq C \Delta t$, and

$$\|v(x) - v(\hat{\mathcal{X}}_x^l(t_{l-1})^{-1})\|_0^2 \leq \int_{\Omega} \left(\int_{t_{l-1}}^{t_l} \frac{d}{dt} v(\hat{\mathcal{X}}_x^l(t)^{-1}) dt \right)^2 dx \leq C \Delta t^2 \|\nabla v\|_0^2.$$

Similarly, we have

$$\|v(\hat{\mathcal{X}}_x^l(t_{l-1})^{-1})\| \leq (1 + C \Delta t) \|v\|_0^2.$$

Then, we deduce

$$\|\eta^{l-1} - \hat{\eta}^{l-1}\|_{-1} \leq C \Delta t \|\eta^{l-1}\|_0. \quad (4.13)$$

By (4.11) and (4.13), we get

$$\left\| \frac{\eta^l - \hat{\eta}^{l-1}}{\Delta t} \right\|_{-1} \leq C h^{k+1} \|u_h^{l-1}\|_{\infty} + C \frac{h^{k+1}}{\sqrt{\Delta t}} \left\| \frac{\partial u}{\partial t} \right\|_{L^2(t_{l-1}, t_l; H^k(\Omega))}. \quad (4.14)$$

Similarly, we obtain

$$\left\| \frac{\xi_h^{l-1} - \hat{\xi}_h^{l-1}}{\Delta t} \right\|_{-1} \leq C \|\xi_h^{l-1}\|_0. \quad (4.15)$$

Setting $v_h = \xi_h^l$ in (4.9), we obtain

$$\begin{aligned} &\frac{\|\xi_h^l\|_0^2 - \|\xi_h^{l-1}\|_0^2}{2\Delta t} + (\nu + \nu_{add}) a(\xi_h^l, \xi_h^l) \\ &\leq \left| \left(\frac{u^l - \bar{u}_h^{l-1}}{\Delta t} - \nu \Delta u^l + \nabla p^l - f^l, \xi_h^l \right) \right| + \left| \left(\frac{\eta^l - \hat{\eta}_h^{l-1}}{\Delta t}, \xi_h^l \right) \right| \\ &\quad + \left| \left(\frac{\hat{u}^{l-1} - \bar{u}^{l-1}}{\Delta t}, \xi_h^l \right) \right| + \left| \left(\frac{\hat{\xi}_h^{l-1} - \xi_h^{l-1}}{\Delta t}, \xi_h^l \right) \right| \\ &\quad + \left| \nu_T \left((I - P_{L_H}) \nabla \eta_h^n, (I - P_{L_H}) \nabla \xi_h^l \right) \right| + \left| \nu_T \left((I - P_{L_H}) \nabla u^l, (I - P_{L_H}) \nabla \xi_h^l \right) \right|. \end{aligned}$$

Using Cauchy-Schwarz and Young's inequalities, we deduce

$$\begin{aligned}
& \frac{\|\xi_h^l\|_0^2 - \|\xi_h^{l-1}\|_0^2}{2\Delta t} + (\nu + \nu_{add})\|\nabla \xi_h^l\|_0^2 \\
& \leq C \left\| \frac{u^l - \bar{u}_h^{l-1}}{\Delta t} - \nu \Delta u^l + \nabla p^l - f^l \right\|_0 \|\nabla \xi_h^l\|_0 + C \left\| \frac{\eta^l - \hat{\eta}_h^{l-1}}{\Delta t} \right\|_{-1} \|\nabla \xi_h^l\|_0 \\
& + C \left\| \frac{\hat{u}^{l-1} - \bar{u}^{l-1}}{\Delta t} \right\|_0 \|\nabla \xi_h^l\|_0 + C \left\| \frac{\hat{\xi}_h^{l-1} - \xi_h^{l-1}}{\Delta t} \right\|_{-1} \|\nabla \xi_h^l\|_0 \\
& + \nu_{add} \|\nabla \eta_h^l\|_0 \|\nabla \xi_h^l\|_0 + \nu_T \|(I - P_{L_H}) \nabla u^l\|_0 \|(I - P_{L_H}) \nabla \xi_h^l\|_0 \\
& \leq C \left(\left\| \frac{u^l - \bar{u}_h^{l-1}}{\Delta t} - \nu \Delta u^l + \nabla p^l - f^l \right\|_0^2 + \left\| \frac{\eta^l - \hat{\eta}_h^{l-1}}{\Delta t} \right\|_{-1}^2 + \left\| \frac{\hat{u}^{l-1} - \bar{u}^{l-1}}{\Delta t} \right\|_0^2 \right. \\
& \left. + \left\| \frac{\hat{\xi}_h^{l-1} - \xi_h^{l-1}}{\Delta t} \right\|_{-1}^2 + \nu_{add} \|\nabla \eta_h^l\|_0^2 + \nu_T \|(I - P_{L_H}) \nabla u^l\|_0^2 \right) + \frac{\nu + \nu_{add}}{2} \|\nabla \xi_h^l\|_0^2.
\end{aligned}$$

By (4.1), (4.10), (4.14) and (4.15), we get

$$\begin{aligned}
& \frac{\|\xi_h^l\|_0^2 - \|\xi_h^{l-1}\|_0^2}{\Delta t} + (\nu + \nu_{add})\|\nabla \xi_h^l\|_0^2 \leq C(\Delta t^2 + h^{2k+2} + \nu_{add}h^{2k} \\
& + \frac{h^{2k+2}}{\Delta t} \left\| \frac{\partial u}{\partial t} \right\|_{L^2(t_{l-1}, t_l; H^k(\Omega))}^2 + \|\xi_h^{l-1}\|_0^2 + \nu_T \|(I - P_{L_H}) \nabla u^l\|_0^2).
\end{aligned} \tag{4.16}$$

Summing inequality (4.16) from $i = 1$ to l , we obtain

$$\begin{aligned}
& \|\xi_h^l\|_0^2 + \nu \Delta t \sum_{i=1}^l \|\nabla \xi_h^i\|_0^2 \\
& \leq C \left(\Delta t^2 + h^{2k+2} + \nu_{add}h^{2k} + \nu_T \|(I - P_{L_H}) \nabla u^l\|_0^2 + \Delta t \sum_{i=1}^{l-1} \|\xi_h^i\|_0^2 \right).
\end{aligned}$$

Using discrete Gronwall's Lemma, we have

$$\begin{aligned}
& \|\xi_h^l\|_0^2 + \nu \Delta t \sum_{i=1}^l \|\nabla \xi_h^i\|_0^2 \\
& \leq C \exp(CT) (\Delta t^2 + h^{2k+2} + \nu_{add}h^{2k} + \nu_T \|(I - P_{L_H}) \nabla u^l\|_0^2). \tag{4.17}
\end{aligned}$$

By triangle inequality, we deduce

$$\|u_h^l\|_\infty \leq \|u_h^l - R_h^l\|_\infty + \|R_h^l\|_\infty.$$

Using inverse inequality, $\|v_h\|_\infty \leq Ch^{-1} \|\nabla v_h\|_0$ (see [4]), we get

$$\begin{aligned}
& \|u_h^l\|_\infty \leq Ch^{-1} \|\nabla(u_h^l - R_h^l)\|_0 + \|R_h^l\|_\infty, \\
& \Delta t \|u_h^{n+1}\|_{W^{1,\infty}} \leq \Delta t \|R_h^{l+1}\|_{W^{1,\infty}} + \Delta t \|\xi_h^{l+1}\|_{W^{1,\infty}} \leq C\Delta t + Ch^{-1} \Delta t \|\xi_h^{l+1}\|_1.
\end{aligned}$$

Using (4.17) and Lemma 4.4 we finish the proof.

Theorem 1. Let u_h^n be defined by (3.1) and u be the solution of (1.1). Under the assumptions of Lemma 4, for all $1 \leq n \leq N$ we have

$$\begin{aligned} \|u^n - u_h^n\|_0^2 &\leq C \exp(CT) (\Delta t^2 + h^{2k+2} + \nu_{add} h^{2k} + \nu_T \|(I - P_{L_H}) \nabla u^n\|_0^2), \\ \nu \Delta t \sum_{i=1}^n \|\nabla(u^n - u_h^n)\|_0^2 &\leq C \exp(CT) \\ &\times (\Delta t^2 + h^{2k+2} + \nu_{add} h^{2k} + \nu_T \|(I - P_{L_H}) \nabla u^n\|_0^2) + Ch^{2k}. \end{aligned} \quad (4.18)$$

Proof. Using triangle inequality, (4.3) and (4.4), we obtain the desired result.

Corollary 1. Under the assumption of Lemma 4, Remark 3.2, the regularity assumption of $(u, p) \in (H^3(\Omega)^2 \cap X, H^2(\Omega) \cap M) \forall t \in (0, T]$ and the assumption of ν_T , for all $1 \leq n \leq N$ we have the following error analysis

$$\begin{aligned} \|u^n - u_h^n\|_0^2 &\leq C \exp(CT) (\Delta t^2 + h^{2(k+1)}), \\ \nu \Delta t \sum_{i=1}^n \|\nabla(u^n - u_h^n)\|_0^2 &\leq C \exp(CT) (\Delta t^2 + h^{2k}), \quad k = 1, 2. \end{aligned}$$

4.2 Error estimate for pressure

The following result on the pressure is a consequence of the previous theorem on the velocity.

Theorem 2. Under the assumptions of Lemma 4, for all $1 \leq n \leq N$ we have

$$\left(\sum_{i=0}^n \Delta t \|p^i - p_h^i\|_0^2 \right)^{1/2} \leq C(\Delta t + h^k + \nu_T \|(I - P_{L_H}) \nabla u_h^n\|_0),$$

where C is a positive constant independent of Δt and h .

Proof. Multiplying (1.1) by $v_h \in X_h$ and subtracting the result from (3.6) with $\varphi_h = 0$, we obtain

$$\begin{aligned} d(Q_h^n - p_h^n, v_h) &= -(\bar{e}(x, n), v_h) + d(Q_h^n - p^n, v_h) + \nu a(u_h^n - u^n, v_h) \\ &+ \left(\frac{(u^n - u_h^n) - (\hat{u}_h^{n-1} - \bar{u}_h^{n-1})}{\Delta t}, v_h \right) + \nu_T \left((I - P_{L_H}) \nabla u_h^n, (I - P_{L_H}) \nabla v_h \right). \end{aligned}$$

By the inf-sup condition and Cauchy-Schwarz inequality, we get

$$\begin{aligned} \|Q_h^n - p_h^n\|_0 &\leq C \left(\|e(x, n)\|_0 + \left\| \frac{u - u_h^n}{\Delta t} \right\|_0 + \left\| \frac{\hat{u}_h^{n-1} - \bar{u}_h^{n-1}}{\Delta t} \right\|_0 + \|Q_h^n - p^n\|_0 \right) \\ &+ \nu \|\nabla(u_h^n - u^n)\|_0 + \nu_T \|(I - P_{L_H}) \nabla u_h^n\|_0. \end{aligned} \quad (4.19)$$

Using (4.1), (4.10) and (4.18), we arrive at

$$\left(\sum_{i=0}^n \Delta t \|Q_h^i - p_h^i\|_0^2 \right)^{1/2} \leq C(\Delta t + h^k + \nu_T \|(I - P_{L_H}) \nabla u_h^n\|_0).$$

By triangle inequality, we have

$$\begin{aligned} \left(\sum_{i=0}^n \Delta t \|p^i - p_h^i\|_0^2\right)^{1/2} &\leq \left(\sum_{i=0}^n \Delta t \|Q_h^i - p_h^i\|_0^2\right)^{1/2} + \left(\sum_{i=0}^n \Delta t \|Q_h^i - p^i\|_0^2\right)^{1/2} \\ &\leq C(\Delta t + h^k + \nu_T \|(I - P_{L_H})\nabla u_h^n\|_0). \end{aligned}$$

Therefore, we finish the proof. \square

Corollary 2. *Under the assumptions of Lemma 4 and Remark 3.2, the regularity assumption of $(u, p) \in (H^3(\Omega)^2 \cap X, H^2(\Omega) \cap M) \forall t \in (0, T]$ and the assumption of ν_T , we have the following error estimate*

$$\left(\sum_{i=0}^n \Delta t \|p^i - p_h^i\|_0^2\right)^{1/2} \leq C(\Delta t + h^k), \quad k = 1, 2. \tag{4.20}$$

5 Numerical Results

In this section, we present some numerical results to verify the theoretical results obtained in the previous section and show the effect of our methods. Here, we use the software package FreeFEM++ [21] for our program.

5.1 Analytical solution problems

In this subsection, we present some numerical results of the Navier-Stokes problems with the analytical solution

$$\begin{aligned} u_1 &= 10x^2(x-1)(2x-1)y^2(y-1)^2 \exp(-2\pi^2 t\nu), \\ u_2 &= -10x(x-1)(2x-1)y^2(y-1)^2 \exp(-2\pi^2 t\nu), \\ p &= 20(2x-1)(2y-1) \exp(-4\pi^2 t\nu). \end{aligned}$$

The boundary and initial conditions in (1.1) are set equal to the analytical solution and f is given by evaluating the momentum equation of the problem (1.1) for the analytical solution. We choose $\Delta t = 0.001$, $T = 1$, $\delta = 0.5$ and $Re = 2000$. We present the numerical results with different h respectively by using Taylor-Hood element, and $L_H = \{\mathbb{L} \in (L^2(\Omega))^{2 \times 2}, \mathbb{L}|_K \in P_{1dc}^{2 \times 2} \forall K \in \mathfrak{S}_h\}$, P_{1dc} means the piecewise linear discontinuous finite element space [16].

Table 1. The numerical results of MMOCVMS at $T = 1$ with $\Delta t = 0.001$, $Re = 2000$, $\nu_T = \delta h^2 \|\nabla u_h\|_0$.

$1/h$	$\frac{\ u-u_h\ _0}{\ u\ _0}$	$\frac{\ \nabla(u-u_h)\ _0}{\ u\ _0}$	$\frac{\ p-p_h\ _0}{\ p\ _0}$	u_{L_2} rate	u_{H_1} rate	P_{L_2} rate	CPU, s
10	2.59204e-3	3.06620e-2	7.74597e-3	—	—	—	178.433
20	3.14163e-4	7.47865e-3	1.93649e-3	3.0445	2.0356	2.0000	727.026
30	9.41846e-5	3.29969e-3	8.60663e-4	2.9711	2.0180	2.0000	1650.72
40	4.30740e-5	1.85122e-3	4.84123e-4	2.7195	2.0091	2.0000	2988.12

Table 2. The numerical results of MMOCFEM mixed finite element method at $T = 1$ with $\Delta t = 0.001$, $Re = 2000$.

$1/h$	$\frac{\ u-u_h\ _0}{\ u\ _0}$	$\frac{\ \nabla(u-u_h)\ _0}{\ u\ _0}$	$\frac{\ p-p_h\ _0}{\ p\ _0}$	u_{L_2} rate	u_{H_1} rate	P_{L_2} rate	CPU, s
10	2.59204e-3	3.06620e-2	7.74597e-3	—	—	—	214.547
20	3.14163e-4	7.47865e-3	1.93649e-3	3.0445	2.0356	2.0000	872.899
30	9.41846e-5	3.29969e-3	8.60663e-4	2.97105	2.01799	2.0000	1987.68
40	4.3074e-5	1.85122e-3	4.84123e-4	2.7195	2.0091	2.0000	3553.67

Table 3. The numerical results of MMOCVMS at $T = 1$ with $h = 1/100$, $Re = 2000$, $\nu_T = \delta h^2 \|\nabla u_h\|_0$.

Δt	$\frac{\ u-u_h\ _0}{\ u\ _0}$	$\frac{\ \nabla(u-u_h)\ _0}{\ u\ _0}$	$\frac{\ p-p_h\ _0}{\ p\ _0}$	u_{L_2} rate	u_{H_1} rate	CPU, s
0.1	1.81239e-3	2.23524e-2	7.74605e-5	—	—	207.386
0.05	9.06667e-4	1.16259e-3	7.74599e-5	1.13258	1.06939	376.132
0.025	4.5537e-4	6.46561e-4	7.74597e-5	0.9935	0.8464	752.062
0.0125	2.28506e-4	4.16179e-4	7.74597e-5	0.9774	0.6272	1525.39

The numerical results in Table 1 and 2 show that the convergence rates are $\mathcal{O}(h^3)$ of the L^2 -norm for u , $\mathcal{O}(h^2)$ of the semi H^1 -norm for u and $\mathcal{O}(h^2)$ of the L^2 -norm for p , which agree very well with our theoretical results by using $P_2 - P_1$ finite element spaces. In Table 3 and 4, we present the convergence test results for time. We can see that the time convergence order is $\mathcal{O}(\Delta t)$, which agrees well with our theory analysis. Comparing our numerical results with the MMOCFEM's, we can see that the error and the convergence rates are very similar. But our method can save much time.

5.2 The lid driven cavity problem

In this subsection we show the numerical results of lid driven cavity problem. The two-dimensional lid driven was formulated as in $\Omega = (0, 1)^2$, the boundary conditions are $u_1 = 1, u_2 = 0$ on the top lid and $u_1 = 0, u_2 = 0$ on the other lids. In a former work [5], it was suggested that the first Hopf bifurcation occurs around Reynolds number $Re = 7500$. Since then various results were given in the literature [5, 12, 14]. In this numerical experiments, $h = \frac{1}{60}$, $\Delta t = 0.01$, $\nu_T = \delta h^2 \|\nabla u_h\|_0$, $\delta = 0.5$ are chosen. We choose the Taylor-Hood element and $L_H = \{\mathbb{L} \in (L^2(\Omega))^{2 \times 2}, \mathbb{L}|_K \in P_{1dc}^{2 \times 2} \forall K \in \mathfrak{S}_h\}$.

Firstly, we choose $Re = 5000$, which is a good choice as there are some comparisons available in the literature and as the steady solution is still stable but not too far from the first Hopf bifurcation. Figure 1 gives the numerical results, (a) the numerical streamline at $t = 250$, (b) the numerical pressure contours at $t = 250$, (c) the computed u -velocity profiles along a vertical line passing through the geometric center of the cavity, (d) the computed v -velocity profiles along a horizontal line passing through the geometric center of the cavity, (e) the evolution of the kinetic energy in time and (f) the evolution of $\|u_h^n - u_h^{n-1}\|_0 / \|u_h^n\|_0$ in time.

Table 4. The numerical results of MMOCFEM at $T = 1$ with $h = 1/100$, $Re = 2000$.

Δt	$\frac{\ u-u_h\ _0}{\ u\ _0}$	$\frac{\ \nabla(u-u_h)\ _0}{\ \nabla u\ _0}$	$\frac{\ p-p_h\ _0}{\ p\ _0}$	u_{L_2} rate	u_{H_1} rate	CPU, s
0.1	1.81239e-3	2.23524e-2	7.74605e-3	—	—	211.911
0.05	9.06667e-4	1.16259e-3	7.74599e-5	1.13258	1.06939	385.82
0.025	4.5537e-4	6.46561e-4	7.74597e-5	0.9935	0.8464	757.553
0.0125	2.28506e-4	4.16179e-4	7.74597e-5	0.9774	0.6272	1536.96

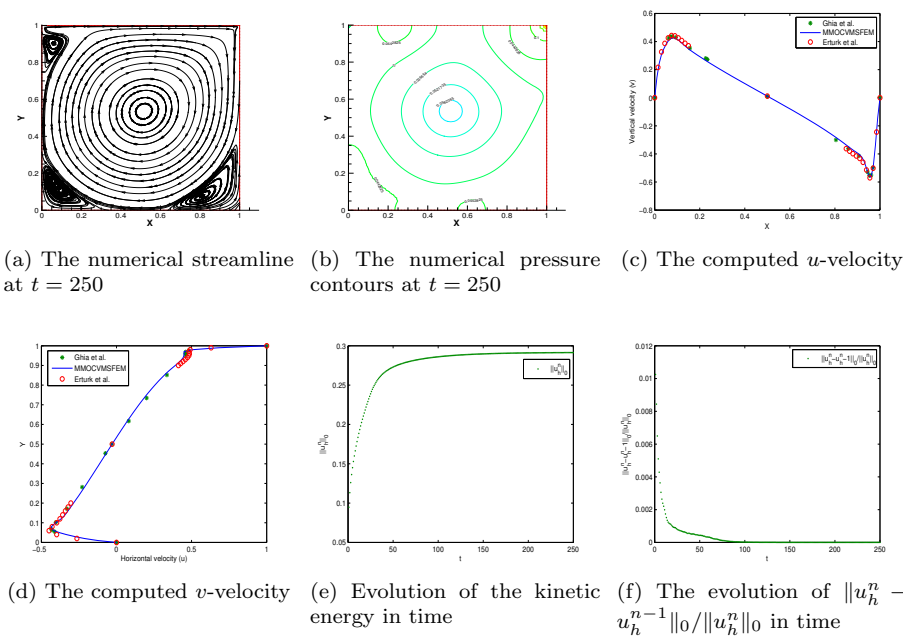


Figure 1. The numerical results for $Re = 5000$ with $\nu_T = \delta h^2 \|\nabla u_h\|_0$.

Figure 2 gives the numerical results for $Re = 7500$, (a) the numerical streamline at $t = 250$, (b) the numerical pressure contours at $t = 250$, (c) the computed u -velocity profiles along a vertical line passing through the geometric center of the cavity, (d) the computed v -velocity profiles along a horizontal line passing through the geometric center of the cavity, (e) the evolution of the kinetic energy in time and (f) the evolution of $\|u_h^n - u_h^{n-1}\|_0 / \|u_h^n\|_0$ in time.

From the numerical results, we can see that the kinetic energy reaches a stable state, $\|u_h^n - u_h^{n-1}\|_0 / \|u_h^n\|_0$ changes very small. It means that we get a steady solution of the time-dependent Navier-Stokes equations. And, comparing with the results of the steady Navier-Stokes equations given by Ghia et al. [14] and Erturk et al. [12], we can see that u -velocity profiles along a vertical line passing through the geometric center of the cavity and v -velocity profiles along a horizontal line passing through the geometric center of the

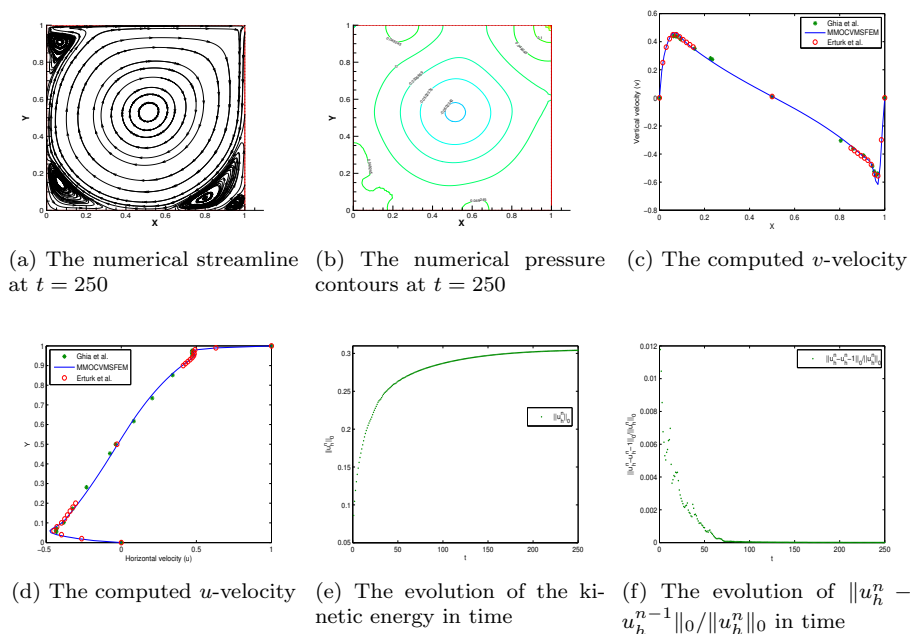


Figure 2. The numerical results for $Re = 7500$ with $\nu_T = \delta h^2 \|\nabla u_h\|_0$.

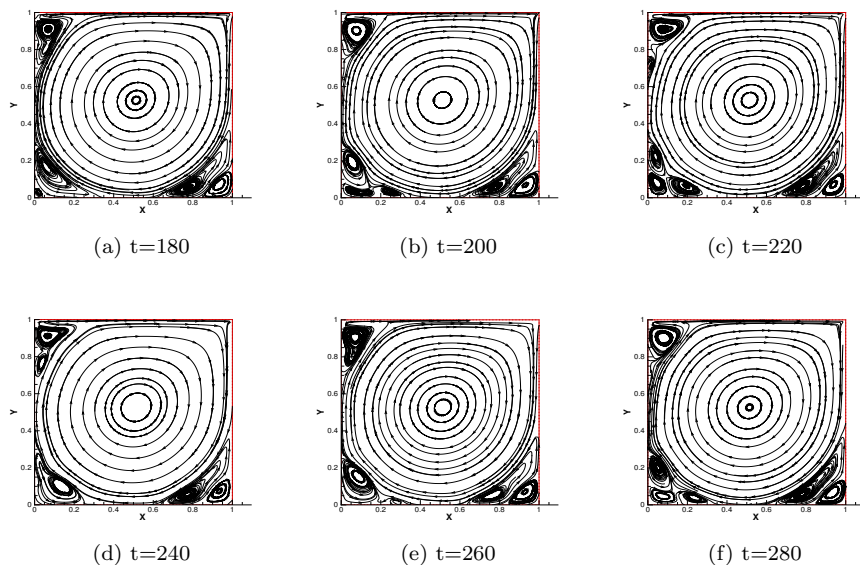


Figure 3. The numerical velocity for $Re = 10000$ with $\nu_T = \delta h^2 \|u_h^n\|_0$ at different time.

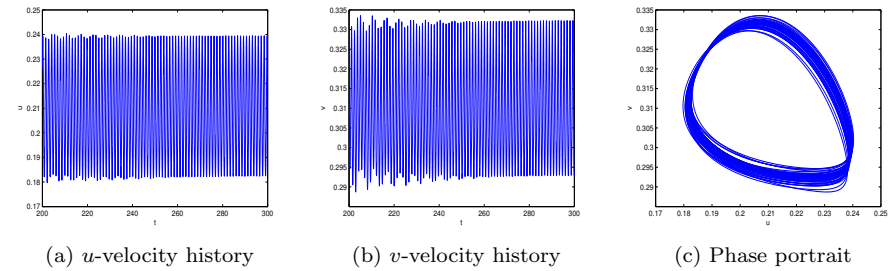


Figure 4. u -velocity history (a), v -velocity history (b) and the phase portrait (c), at monitoring point $(1/8, 13/16)$ for $Re = 10000$ with $\nu_T = \delta h^2 \|u_h^n\|_0$ from $t = 200$ to $t = 300$.

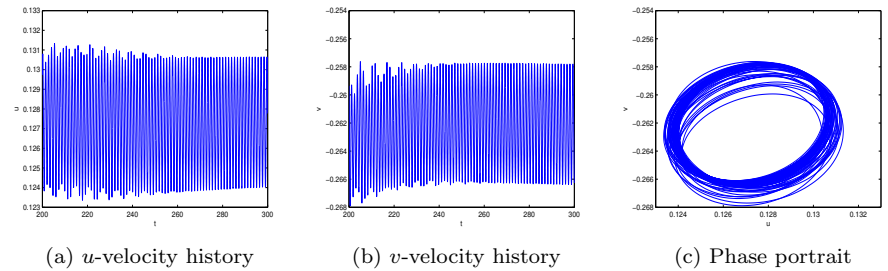


Figure 5. u -velocity history (a), u -velocity history (b) and the phase portrait (c), at monitoring point $(7/8, 13/16)$ for $Re = 10000$ with $\nu_T = \delta h^2 \|u_h^n\|_0$ from $t = 200$ to $t = 300$.

cavity agree very well.

Finally, we show some numerical results for $Re = 10000$, which is probably the most famous value and for quite a long time the question was to know if the steady solution was stable or not for this Reynolds number. In [6], Bruneau and Saad shown that the steady is not stable and the first Hopf bifurcation occurs around $Re = 8000$. Figure 3 presents the streamlines at different time. We choose two monitoring points $(1/8, 13/16)$ and $(7/8, 13/16)$. Figures 4 and 5 show u -velocity history (a), v -velocity history (b) and the phase (c) portrait at monitoring points $(1/8, 13/16)$ and $(7/8, 13/16)$. It shows that the stable solution in mainly periodic with small variations in the amplitude of the time evolution at the monitoring points. And the phase portraits show that the variations in amplitude yield a solution which is quasi-periodic. Figure 6 shows the evolution of the kinetic energy in time (a) and evolution of $\|u_h^n - u_h^{n-1}\|_0 / \|u_h^n\|_0$ in time (b). We can see that the kinetic energy dose not change as the time change long enough and the errors don't change smaller. The results are very close the those shown in [6].

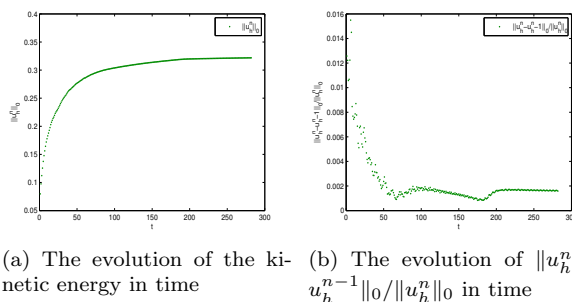


Figure 6. Evolution of the kinetic energy in time (a) and evolution of $\|u_h^n - u_h^{n-1}\|_0 / \|u_h^n\|_0$ in time (b) for $Re = 10000$ with $\nu_T = \delta h^2 \|u_h^n\|_0$.

5.3 The flow around a cylinder problems

In this part, we investigate the two-dimensional under-resolved fluid flow around a cylinder [37], the physical model given by Figure 7. From Figure 7 we can see that, the height of the channel $H = 0.41$ and the width $W = 2.2$. The origin of the cylinder is at $(0.2, 0.2)$, and the radius R is equal to 0.05 . The time-dependent inlet flow velocity profiles are given by

$$u_1 = 6y(1 - H)/H^2, \quad u_2 = 0.$$

The outlet boundary condition is given by

$$\frac{\partial u}{\partial n} = 0.$$

The boundary conditions of the two parallel planes and the cylinder surface are set as the non-slip boundary conditions. The Reynolds number is defined by $Re = 2\bar{U}R/\nu$ and $\bar{U} = 2u(0, H/2)/3$. Here, we choose $Re = 100$, $\Delta t = 0.01$. The finite element grid is presented by Figure 7. We choose the Taylor-Hood element spaces. Figures 8 and 9 give the numerical results by the MMOCVMSFEM. Figures 10 and 11 present the numerical results by the MMOCFEM. From the numerical results, we can see that the numerical results of the MMOCVMSFEM agree well with the numerical results of MMOCFEM.

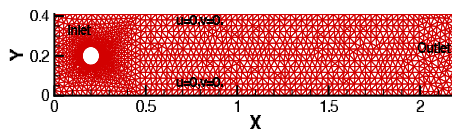


Figure 7. Physical model of the flow around a cylinder problems.

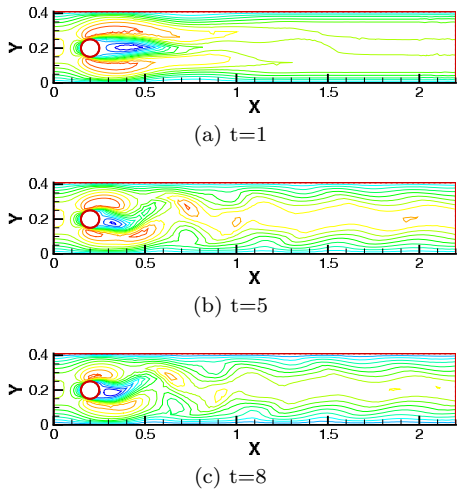


Figure 8. The numerical contours of the velocity u at different time by MMOCVMSFEM.

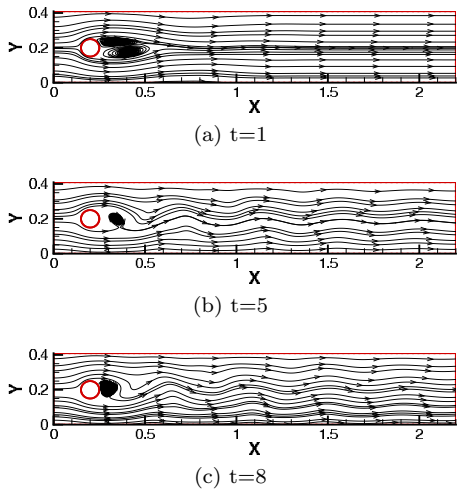


Figure 9. The numerical streamlines at different time by MMOCVMSFEM.

6 Conclusions

In this work, we proposed a modified algorithm of characteristics variational multiscale finite element method for the time dependent Navier-Stokes equations. The theoretical analysis is discussed and illustrated. Error estimate shows that this method has a good convergence property. Numerical experiments are presented to demonstrate the high efficiency of the new algorithm. Also, further developments can extend these techniques and ideas to the general nonlinear problems and high-dimensional problems.

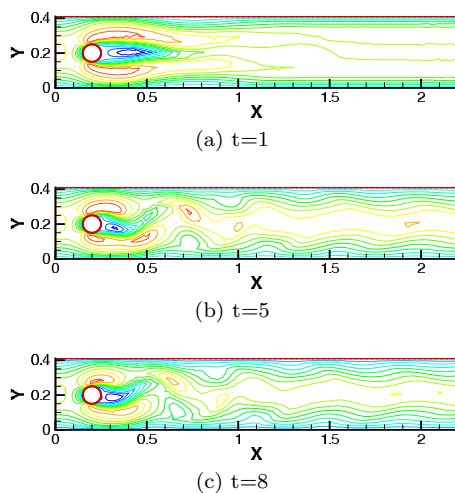


Figure 10. The numerical contours of the velocity u at different time by MMOCFEM.

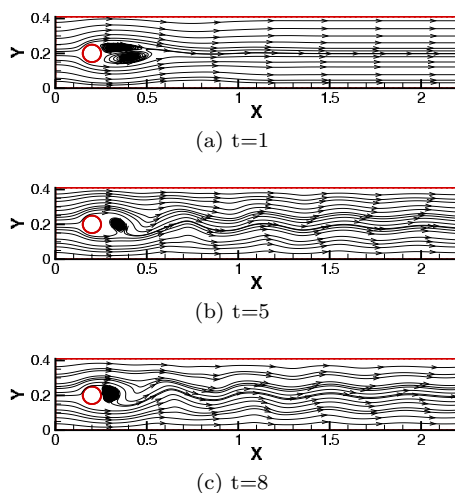


Figure 11. The numerical streamlines at different time by MMOCFEM.

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